

## Supplement to Lec 17

We wish to give a "slicker" proof of the relation

$$\begin{aligned}\xi(z) &= \xi(1-z), \\ \xi(z) &\equiv \pi^{-z/2} \Gamma\left(\frac{z}{2}\right) J(z).\end{aligned}$$

We use the Poisson summation formula.

Consider functions  $f \in C^\infty(\mathbb{R})$  such that

$$|f^{(j)}(x)| \leq A_j \Delta (1+|x|)^{-\Delta}$$

for every  $j \geq 0$  and  $\Delta \geq 1$ . This is the so-called Schwartz class of functions on  $\mathbb{R}$ .

Let

$$\hat{f}(u) = \int_{\mathbb{R}} f(x) e^{-2\pi i u x} dx$$

← the Fourier transform of  $f$

By repeatedly integrating by parts,

$$|\hat{f}(u)| \leq B_\Delta (1+|u|)^{-\Delta}, \quad u \in \mathbb{R}$$

for every  $\Delta \geq 1$ .

(2)

A similar inequality holds for any derivative  $\hat{f}^{(m)}(u)$  after using Leibnitz's rule for improper integrals.

Thus,  $\hat{f}^1$  is also in the Schwartz class of functions on  $\mathbb{R}$ .

The function

$$g(x) \equiv \sum_{n=-\infty}^{\infty} f(x+n), \quad x \in \mathbb{R}$$

is very interesting. \*

\* Remember that an infinite series  $\sum_{k=1}^{\infty} u_k(x)$  can be differentiated term-by-term when  $u_k \in C^1(\mathbb{R})$ ,  $\sum_1^{\infty} u_k'(x)$  conv unif on  $\mathbb{R}$  compacta, and  $\sum_1^{\infty} u_k(x)$  conv pointwise.

Clearly RHS converges unif + absolutely <sup>(3)</sup> on  $\mathbb{R}$ -compacta. We also have:

$$g(x+1) = g(x) ;$$

$$g^{(j)}(x) = \sum_{n=-\infty}^{\infty} f^{(j)}(x+n) .$$

Here  $j \geq 1$  and the RHS again converges unif + absolutely on  $\mathbb{R}$  compacta.

The function  $g$  is thus  $C^\infty$  and periodic. Expand it as a Fourier series !!!

$$g(x) = \sum_{k=-\infty}^{\infty} C_k e^{2\pi i k x}$$

$$C_k = \int_0^1 g(x) e^{-2\pi i k x} dx$$

à la  
Lec. 12  
p. 18

But,

$$C_k = \lim_{N \rightarrow \infty} \int_0^1 \left( \sum_{-N}^N f(x+n) \right) e^{-2\pi i k x} dx$$

$$= \lim_{N \rightarrow \infty} \sum_{-N}^N \int_n^{n+1} f(y) e^{-2\pi i k y} dy$$

$$= \lim_{N \rightarrow \infty} \int_{-N}^{N+1} f(y) e^{-2\pi i k y} dy \quad (4)$$

$$= \int_{-\infty}^{\infty} f(y) e^{-2\pi i k y} dy = \hat{f}(k) .$$

In other words:  $(x \in \mathbb{R})$

$$\sum_{n=-\infty}^{\infty} f(x+n) = \sum_{k=-\infty}^{\infty} \hat{f}(k) e^{2\pi i k x} .$$

Both sides are unif + abs convergent on  $\mathbb{R}$ -compacta.

Put  $x = 0$ . Get:

$$\boxed{\sum_{n=-\infty}^{\infty} f(n) = \sum_{k=-\infty}^{\infty} \hat{f}(k)}$$

This is the famous Poisson summation formula.   
 {NB: The conditions on  $f(x)$  can be relaxed significantly.}

One knows that

$$\int_{-\infty}^{\infty} e^{-\varphi x^2} dx = \sqrt{\frac{\pi}{\varphi}} \quad (\varphi > 0) .$$

By baby complex analysis,

← CIT (5)

$$f(x) = e^{-qx^2}$$

$$\Rightarrow \hat{f}(u) = \sqrt{\frac{\pi}{q}} e^{-\frac{\pi u^2}{q}}$$

Thus,

$$\sum_{n=-\infty}^{\infty} e^{-qn^2} = \sqrt{\frac{\pi}{q}} \sum_{k=-\infty}^{\infty} e^{-\frac{\pi k^2}{q}}$$

Put  $q = \pi\beta$  to rephrase this as:

$$\sum_{n=-\infty}^{\infty} e^{-\pi n^2 \beta} = \sqrt{\frac{1}{\beta}} \sum_{k=-\infty}^{\infty} e^{-\frac{\pi}{\beta} k^2}, \quad \beta > 0.$$

(This is the so-called THETA IDENTITY.)

With

$$\Psi(x) \equiv \sum_{n=1}^{\infty} e^{-\pi n^2 x} \quad (x > 0),$$

we have

$$1 + 2\Psi(x) = \sqrt{\frac{1}{x}} [1 + 2\Psi(\frac{1}{x})]$$

Riemann found a nice formula for  $\zeta(s)$  (6)  
using  $\Psi(x)$ . It gives  $\zeta(s) = \zeta(1-s)$  immediately.

We follow Riemann and put  
 $s \equiv \sigma + it$ .

Let  $\mathcal{R} = \{ \operatorname{Re}(s) > 1 \}$ . Let  $K$  be any closed  
disk in  $\mathcal{R}$ . Keep  $s \in K$ .

Note that: p. (5)

$\Psi \in C[(0, \infty)]$

$$\Psi(x) = O(e^{-\pi x}) \quad , \quad x \rightarrow +\infty$$

$$\Psi(x) = O\left(\frac{1}{\sqrt{x}}\right) \quad , \quad x \rightarrow 0^+$$

Look at:  $\int_0^\infty \Psi(x) x^{\frac{s}{2}-1} dx$  (principal value)

$$\int_0^\infty \Psi(x) x^{\frac{s}{2}-1} dx \quad (s \in K).$$

This integral satisfies a nice Weierstrass  
M-test for  $s \in K$ . We propose to  
integrate term-by-term.

At once:

$$\begin{aligned} \text{integral} &= \sum_{n=1}^{\infty} \left( \int_0^{\infty} e^{-\pi n^2 x} x^{s/2} \frac{dx}{x} \right) \\ &= \sum_{n=1}^{\infty} \left( \int_0^{\infty} e^{-y} \frac{y^{s/2}}{\pi^{s/2} n^s} \frac{dy}{y} \right) \\ &= \pi^{-s/2} \sum_{n=1}^{\infty} n^{-s} \left( \int_0^{\infty} y^{\frac{s}{2}-1} e^{-y} dy \right) \\ &= \pi^{-s/2} \Gamma\left(\frac{s}{2}\right) \zeta(s) \end{aligned}$$

$$x = \frac{y}{\pi n^2}$$

NICE!

On the other hand,

$$\begin{aligned} \text{integral} &= \int_0^1 \Psi(x) x^{\frac{s}{2}-1} dx \\ &\quad + \int_1^{\infty} \Psi(x) x^{\frac{s}{2}-1} dx \end{aligned}$$

SEK

$$\left\{ \begin{array}{l} \text{but} \\ \Psi(x) = -\frac{1}{2} + \frac{1}{2\sqrt{x}} + \frac{1}{\sqrt{x}} \Psi\left(\frac{1}{x}\right) \end{array} \right\}$$

$$\begin{aligned} &= \int_0^1 \left[ -\frac{1}{2} + \frac{1}{2\sqrt{x}} + \frac{1}{\sqrt{x}} \Psi\left(\frac{1}{x}\right) \right] x^{\frac{s}{2}-1} dx \\ &\quad + \int_1^{\infty} \Psi(x) x^{\frac{s}{2}-1} dx \end{aligned}$$

$$\begin{aligned}
&= -\frac{1}{2} \int_0^1 x^{\frac{s}{2}-1} dx + \frac{1}{2} \int_0^1 x^{\frac{s}{2}-\frac{3}{2}} dx \\
&\quad + \int_0^1 \Psi\left(\frac{1}{x}\right) x^{\frac{s}{2}-\frac{3}{2}} dx + \int_1^\infty \Psi(x) x^{\frac{s}{2}-1} dx \\
&\approx -\frac{1}{s} + \frac{1}{s-1} + \int_1^\infty \Psi(y) y^{\frac{s}{2}-\frac{3}{2}} \frac{dy}{y^2} \\
&\quad + \int_1^\infty \Psi(x) x^{\frac{s}{2}-1} dx \\
&\approx -\frac{1}{s} + \frac{1}{s-1} + \int_1^\infty \Psi(x) \left[ x^{\frac{s}{2}-1} + x^{-\frac{1}{2}-\frac{s}{2}} \right] dx \\
&= \frac{1}{s(s-1)} + \int_1^\infty \Psi(x) \left[ x^{\frac{s}{2}} + x^{\frac{1-s}{2}} \right] \frac{dx}{x}
\end{aligned}$$

continuous and  $O(e^{-\pi x})$  for  $x$  large

In other words,

$$(*) \quad \pi^{-\frac{s}{2}} \Gamma\left(\frac{s}{2}\right) \zeta(s) = \frac{1}{s(s-1)} + \int_1^\infty \Psi(x) \left[ x^{\frac{s}{2}} + x^{\frac{1-s}{2}} \right] \frac{dx}{x}$$

for every  $s \in K$ . NOTICE THAT THE RHS is a perfectly good analytic function of  $s$  on  $\mathbb{C} - \{0, 1\}$ ; the integral is actually an entire function.

THINK ANALYTIC !!  
CONTINUATION  $\infty$

Moreover, ~~the~~ <sup>RHS</sup> remains the same when

$$s \longleftrightarrow 1-s$$

QED ■