# INHOMOGENEOUS RANDOM GRAPHS

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## 1. Background

One of the most studied random graphs is $G(n,p)$, which has $n$ vertices that can be taken as the integers $1,\dots,n$, and where each pair of vertices is connected by an edge with probability $p$, independently of all other edges.

Consider the case $p = c/n$ for some constant $c > 0$, and let $n \to \infty$. The degree of a given vertex has a binomial distribution $\mathrm{Bi}(n-1, c/n) \approx \mathrm{Po}(c)$. This is a strongly concentrated distribution with an exponentially decreasing tail. Many graphs from "real life" have degree distributions with much larger tails, for example power-law tails. One popular example is the Internet, either the physical network of servers or the (directed) network of web pages and links. Other examples that have been studied include, for example, graphs describing telephone calls, scientific collaborations, and protein interactions.

It is therefore important to study also random graph models with less homogeneity, and in particular models with larger tails for the degree distribution. Indeed, several such models have been proposed and studied, more or less rigorously, see for example Molloy and Reed [36; 37], Kleinberg, Kumar, Raghavan, Rajagopalan and Tomkins [33], Barabási and Albert [4], Aiello, Chung and Lu [1], Bollobás, Riordan, Spencer and Tusnády [12], Cooper and Frieze [19], Bollobás and Riordan [10], Buckley and Osthus [16] and the surveys by Albert and Barabási [2], Dorogovtsev and Mendes [22] and Bollobás and Riordan [8].

We will describe one class of random graphs that generalize $G(n,c/n)$ but also allow many less homogeneous examples, for example natural examples of 'scale-free' random graphs, where the degree distribution has a power-law tail. We believe that when it comes to modelling real-world graphs with, for example, observed power-laws for vertex degrees, our model provides an interesting and flexible alternative to existing models.

Nevertheless, we will see that many properties of $G(n,c/n)$ extend to these random graphs. In particular, we consider the question whether there exist a giant component or not, and we will, typically, find a phase transition similar to what happens for $G(n,c/n)$. There are, however, some interesting twists for some examples.

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We are interested in graphs with a large number of vertices, and in particular in asymptotics as the number tends to infinity. The graphs we consider are such that the average degree stays bounded, so they are rather sparse.

These lectures are based on Bollobás, Janson and Riordan [6]. See this paper for proofs and further details, as well as for further references.

## 2. Some examples

We will study random graphs where the edges appear independently, as in $G(n,p)$, but where the probability of an edge may differ between different pairs of vertices. We begin with some examples of the type of random graph we consider; further examples are given in Sections 6 and 15–21.

**Example 2.1.** Let, as above, the vertices be the $n$ integers $1,\dots,n$. For every pair $\{i,j\}$ with $i\neq j$, independently of all other pairs, connect $i$ and $j$ by an edge with probability

$$p_{ij}=\frac{\kappa(i/n,j/n)}{n} \tag{2.1}$$

for a given symmetric non-negative function $\kappa$ on $(0,1]^2$. Taking $\kappa=c$ constant, we obtain $G(n,c/n)$, but other functions $\kappa$ give many other interesting random graphs.

To be precise, this definition requires that $\kappa(i/n,j/n)\le n$ for all $i$ and $j$, since otherwise $p_{ij}>1$. There is no problem if, say, $\kappa$ is bounded and $n$ is large; however, we want to allow unbounded $\kappa$ too. In general, we therefore modify (2.1) to

$$p_{ij}=\min\Bigl(\frac{\kappa(i/n,j/n)}{n},1\Bigr). \tag{2.2}$$

For notational simplicity, we will sometimes ignore this in the formulas below; hence, a probability $p_{ij}>1$ should always be interpreted as 1.

**Example 2.2.** A specific interesting example of the type in Example 2.1 is given by the choice $\kappa(x,y)=c/\max\{x,y\}$; we then connect $i$ and $j$ by an edge with probability

$$p_{ij}=\frac{c}{\max\{i,j\}}.$$

Here $c>0$ is a parameter that will be kept constant as $n$ varies. (We assume for simplicity that $c\le 2$; otherwise we have to truncate $p_{ij}$ at 1 as in (2.2).) This example is the *uniformly grown random graph*, or *$c/j$-graph*, $G_n^{1/j}(c)$. The graph $G_n^{1/j}(c)$ is thus the graph on $\{1,2,\dots,n\}$ in which edges are present independently, and the probability that for $i\neq j$ the edge $ij$ is present is $p_{ij}=c/\max\{i,j\}$, or simply $c/j$ if $i<j$. We return to this example in Section 16.

In principle, we may take any symmetric non-negative function $\kappa$ on $(0,1]^2$ in Example 2.1. However, the function $\kappa$ is evaluated only at rational points, so in order to make sense of having $\kappa$ defined on the entire square, it is reasonable to impose a continuity condition.

To assume that $\kappa$ is continuous (and thus bounded) on the closed unit square $[0,1]^2$ would be convenient, but too strong for our purposes since it excludes the example $\kappa(x,y) = c/\max\{x,y\}$ just given.

To assume that $\kappa$ is continuous on the open unit square $(0,1)^2$ is enough, but it turns out that it suffices to assume that $\kappa$ is *continuous almost everywhere* in the unit square. This extension is convenient because it allows piecewise continuous functions as in the following simple example.

**Example 2.3.** Let $M \geq 2$ be a fixed integer and divide $(0,1]$ into the $M$ intervals $I_k = ((k-1)/M, k/M]$, $k = 1, \dots, M$. Let $\kappa$ be constant on each square $I_k \times I_l$.

This means that the vertices are of $M$ different types, and that the probability of an edge $ij$ depends (only) on the types of $i$ and $j$. If $n$ is a multiple of $M$, there are $n/M$ vertices of each type. The random graphs obtained in this way are essentially the same as those defined by Söderberg [41; 42; 43; 44], see Example 6.2.

## 3. Variations

We have obtained our edge probabilities $p_{ij}$ by evaluating $\kappa$ at the point $(i/n, j/n)$. An interesting alternative is to let $x_1, \dots, x_n$ be $n$ *random* points in $(0,1]$, independent and uniformly distributed, and then take $p_{ij} = \kappa(x_i, x_j)/n$. (As in (2.1) we divide by $n$ in order to keep the average degree bounded.)

If the vertex labels $1, \dots, n$ do not have any special significance, we may order the sequence $x_1, \dots, x_n$. The ordered sequence is then close to $\frac{1}{n}, \frac{2}{n}, \dots, 1$, so it is not surprising that we will obtain a random graph with the same asymptotic properties as in Example 2.1 above.

As we will see in the general definition below, we can also allow $x_1, \dots, x_n$ to be random and dependent; asymptotically, only the density of the points matter.

There is nothing special with the interval $(0,1]$ here; it can be replaced by another space. For example, the finite-type case in Example 2.3 is simpler described by using a finite type space $\{1, \dots, M\}$, see Example 6.2.

This leads to the general (but long and somewhat technical) definition in the following section.

## 4. Definition

The general inhomogeneous random graph $G^{\mathcal{V}}(n, \kappa)$ is defined as follows. We proceed in two steps, constructing first the vertices and then the edges. Note that $n$ is a parameter measuring the size of the graph, and we are primarily interested in asymptotics as $n \to \infty$.

In many instances, $n$ is the number of vertices, and for simplicity we begin with this case.

A *vertex space* $\mathcal{V}$ is a triple $(\mathcal{S}, \mu, (\mathbf{x}_n)_{n\geq 1})$, such that the following holds.

(i) $\mathcal{S}$ is a separable metric space.

(ii) $\mu$ is a Borel probability measure on $\mathcal{S}$, i.e., a (positive) Borel measure with $\mu(\mathcal{S}) = 1$.
(iii) For each $n$, $\mathbf{x}_n$ is a deterministic or random sequence $(x_1, x_2, \dots, x_n)$ of $n$ points in $\mathcal{S}$.

(Formally, we should write $\mathbf{x}_n = (x_1^{(n)}, \dots, x_n^{(n)})$, say, as we assume no relationship between the elements of $\mathbf{x}_n$ for different $n$, but we omit this extra index.)

In general, however, the number of vertices may be random, and we require only that it is roughly proportional to $n$. We therefore consider also the following extension, using $v_n$ to denote the number of vertices. There is in general no need for the parameter $n$ to be integer valued in this version (but we keep the notation $n$ for consistency), and we assume that $I$ is a given unbounded subset of $(0, \infty)$. Typically, either $I = \mathbb{N}$ or $I = (0, \infty)$.

A *generalized vertex space* $\mathcal{V}$ is a triple $(\mathcal{S}, \mu, (\mathbf{x}_n)_{n\ge1})$, such that the following holds.

(i) $\mathcal{S}$ is a separable metric space.
(ii) $\mu$ is a (positive) Borel measure on $\mathcal{S}$ with $0 < \mu(\mathcal{S}) < \infty$.
(iii) For each $n \in I$, $\mathbf{x}_n$ is a deterministic or random sequence $(x_1, x_2, \dots, x_{v_n})$ of $v_n$ points of $\mathcal{S}$, where $v_n$ may be random.

(Again, we really should write $\mathbf{x}_n = (x_1^{(n)}, \dots, x_{v_n}^{(n)})$, say, but we omit this extra index for simplicity.)

Let $M(\mathcal{S})$ be the space of all (positive) finite Borel measures on $\mathcal{S}$, and equip $M(\mathcal{S})$ with the standard weak topology: $\nu_n \to \nu$ iff $\int f\,d\nu_n \to \int f\,d\nu$ for all bounded continuous functions $f : \mathcal{S} \to \mathbb{R}$. Let

$$\mu_n := \frac{1}{n}\sum_{i=1}^{v_n} \delta_{x_i}$$

where $\delta_x$ is the Dirac measure at $x \in \mathcal{S}$; thus $\mu_n$ is a random element of $M(\mathcal{S})$. Note that the total mass $\mu_n(\mathcal{S}) = v_n/n$; hence $\mu_n$ is a probability measure if and only if $v_n = n$.

We will further assume, both for vertex spaces and generalized vertex spaces, that

(iv) $\mu_n \xrightarrow{\mathrm{p}} \mu$, as elements of $M(\mathcal{S})$.

Recall that a set $A \subseteq \mathcal{S}$ is a *$\mu$-continuity set* if $A$ is (Borel) measurable and $\mu(\partial A) = 0$, where $\partial A$ is the boundary of $A$. The convergence condition (iv) is equivalent to the condition that for every $\mu$-continuity set $A$,

$$\mu_n(A) := \#\{i \le v_n : x_i \in A\}/n \xrightarrow{\mathrm{p}} \mu(A).$$

In particular, it is a consequence of (iv) that $v_n/n \xrightarrow{\mathrm{p}} \mu(\mathcal{S})$. For vertex spaces this holds trivially; for generalized vertex spaces it says that the number of vertices is roughly $\mu(\mathcal{S})n$. (The total mass $\mu(\mathcal{S})$ appears as a simple scale factor; it is convenient for applications to allow general finite values of $\mu(\mathcal{S})$,

but it is easy to reduce to the case $\mu(\mathcal{S}) = 1$ by reparametrizing and changing $n$ to $\mu(\mathcal{S})n$.)

We now complete the construction of $G^{\mathcal{V}}(n,\kappa)$. Let $\mathcal{V}$ be a (generalized) vertex space, and let $\kappa$ be a symmetric non-negative (Borel) measurable function on $\mathcal{S}\times\mathcal{S}$. (We call such a function a *kernel*.) We define the random graph $G^{\mathcal{V}}(n,\kappa)$ by first letting the vertex set be $\{1,\dots,v_n\}$. (It is sometimes more convenient to identify the vertices with the points $x_1,\dots,x_{v_n}$ in $\mathcal{S}$ rather than integers, but note that this must be done with care if there are repetitions among the points $x_i$.)

We then add edges as follows. Given the sequence $\mathbf{x}_n$, we consider each pair of vertices $\{i,j\}$ with $i\neq j$ separately, and let there be an edge between $i$ and $j$ with probability

$$p_{ij} = \min\Big\{\frac{\kappa(x_i,x_j)}{n}, 1\Big\}.$$

This random choice is done independently for all pairs $\{i,j\}$, conditioned on $\mathbf{x}_n$.

In order to avoid pathologies, we finally assume

(v) $\kappa$ is continuous a.e. on $\mathcal{S}\times\mathcal{S}$;

(vi) $\kappa\in L^1(\mathcal{S}\times\mathcal{S},\mu\times\mu)$, i.e.,

$$\iint_{\mathcal{S}^2}\kappa(x,y)\,d\mu(x)\,d\mu(y)<\infty;$$

(vii)

$$\frac{1}{n}\,\mathbb{E}\,e\big(G^{\mathcal{V}}(n,\kappa)\big)\to\frac{1}{2}\iint_{\mathcal{S}^2}\kappa(x,y)\,d\mu(x)\,d\mu(y).$$

We say that the kernel $\kappa$ is *graphical* on the (generalized) vertex space $\mathcal{V}=(\mathcal{S},\mu,(\mathbf{x}_n)_{n\ge1})$ when the conditions (v)–(vii) hold.

It can be shown that (vii) follows from the other assumptions if $\kappa$ is bounded and $\mathcal{V}$ is a vertex space.

As remarked above, the number of vertices $v(G^{\mathcal{V}}(n,\kappa)) = v_n$ is roughly proportional to $n$; more precisely,

$$\frac{v(G^{\mathcal{V}}(n,\kappa))}{n}=\frac{v_n}{n}=\mu_n(\mathcal{S})\xrightarrow{\mathrm{p}}\mu(\mathcal{S}).$$

We sometimes suppress the dependence on $\mathcal{V}$, writing $G(n,\kappa)$ for $G^{\mathcal{V}}(n,\kappa)$.

**Remark 4.1.** Some continuity condition on $\kappa$ is necessary since our definition of vertex spaces is so general, and the a.s. continuity assumed in (v) is a weak restriction. See further the discussion in Section 2.

However, this continuity condition is not necessary in the important case of a vertex space or generalized vertex space where the points $x_i$ are i.i.d., with the common distribution $\mu$ (see Section 3 and Examples 6.4 and 6.5; such vertex spaces are also commonly used in, e.g., the rank one case discussed in Section 19). In fact, in this case, the results can be extended to random graphs defined using a measure space $(\mathcal{S},\mu)$ without any topology

involved. This will be discussed in a future paper Bollobás, Janson and Riordan [7].

**Remark 4.2.** Some condition of the type (vii) is necessary in order to avoid certain pathologies, see Bollobás, Janson and Riordan [6, Example 8.6]. The version given in (vii) is convenient for many applications, since the expectation on the left hand side usually is easily computed, and the condition thus is easy to verify. For example, note that the condition always holds, by a trivial calculation, for a vertex space where the $x_i$ are random i.i.d., with the distribution $\mu$.

Nevertheless, there are applications where the condition is less natural, for example for a generalized vertex space where the number $v_n$ of vertices is random with so large tails that the mean $\mathbb{E}\, v_n = \infty$ (but still $v_n/n \xrightarrow{\mathrm{p}} \mu(\mathcal{S}) < \infty$). In such cases, it is better to replace (vii) by a condition involving convergence in probability rather than convergence of the mean, but we omit the details. (Alternatively, one can reduce to the case above by suitable conditioning, thus eliminating bad events of small probability. In the case with a random $v_n$ just mentioned, one might, for example, condition on $v_n \le 2\mu(\mathcal{S})n$.)

More generally, we may consider a sequence $(\kappa_n)$ of kernels on $\mathcal{S}$ and the corresponding random graphs $G^{\mathcal{V}}(n,\kappa_n)$. We say that the sequence $(\kappa_n)$ of kernels on $(\mathcal{S},\mu)$ is *graphical on $\mathcal{V}$ with limit $\kappa$* if

- (v) $\kappa$ is continuous a.e. on $\mathcal{S}\times\mathcal{S}$;
- (v') for a.e. $(y,z)\in\mathcal{S}^2$, $y_n\to y$ and $z_n\to z$ imply that $\kappa_n(y_n,z_n)\to\kappa(y,z)$;
- (vi) $\kappa\in L^1(\mathcal{S}\times\mathcal{S},\mu\times\mu)$;
- (vii')
$$\frac{1}{n}\,\mathbb{E}\, e(G^{\mathcal{V}}(n,\kappa_n)) \to \tfrac{1}{2}\iint_{\mathcal{S}^2}\kappa(x,y)\,d\mu(x)\,d\mu(y).$$

We will see that most results for $G^{\mathcal{V}}(n,\kappa)$ hold for $G^{\mathcal{V}}(n,\kappa_n)$ too.

## 5. Further variations

A common variation of the construction above is to let the edges be generated with probability

$$p_{ij} := 1-\exp\bigl(-\kappa(x_i,x_j)/n\bigr), \tag{5.1}$$

rather than by $\min\bigl\{\kappa(x_i,x_j)/n,1\bigr\}$ as in (2.2). This is the result if we regard $\kappa$ as intensities of Poisson processes of edges, and construct a multigraph by adding a Poisson number of edges between $i$ and $j$, with mean $\kappa(x_i,x_j)/n$, and then merge multiple edges. (There will be very few multiple edges, and for many purposes it does not matter whether we merge them or not.)

We can treat this version by regarding it as an instance of $G^{\mathcal{V}}(n,\kappa_n)$ with

$$\kappa_n(x,y) := n\bigl(1-\exp(-\kappa(x,y)/n)\bigr).$$

It is easily seen that the sequence $(\kappa_n)$ is graphical with limit $\kappa$ if $\kappa$ is graphical.

Another alternative, studied by Britton, Deijfen and Martin-Löf [14] in a special case, is to let

$$\frac{p_{ij}}{1-p_{ij}} = \frac{\kappa(x_i,x_j)}{n},$$

i.e., to take

$$p_{ij} := \frac{\kappa(x_i,x_j)}{n+\kappa(x_i,x_j)}. \tag{5.2}$$

Again, this is an instance of $G^{\mathcal{V}}(n,\kappa_n)$, now with

$$\kappa_n(x,y) := \frac{\kappa(x,y)}{n+\kappa(x,y)},$$

and again the sequence $(\kappa_n)$ is graphical with limit $\kappa$ if $\kappa$ is graphical.

The version (5.2) is sometimes simpler than our standard one (2.2). In particular, if the sequence $\mathbf{x}_n$ is deterministic, and the kernel $\kappa$ has rank 1, i.e., $\kappa(x,y) = \psi(x)\psi(y)$ for some function $\psi$, then the probability of obtaining a specific graph $G$ on the given vertex set equals

$$C\prod_i \psi(x_i)^{d_i}$$

where $d_i$ is the degree of vertex $i$ and $C$ is a normalization constant. The probability thus depends on the degree sequence only. (In other words, using statistical terminology, the degree sequence is a sufficient statistic.)

Note further that both (5.1) and (5.2) yield $p_{ij} \in [0,1)$, avoiding the minor complication with values greater than 1 that can occur for (2.1).

## 6. More examples

**Example 6.1** (the Erdős-Rényi random graph)**.** If $\kappa = c$ is constant, then the edge probabilities $p_{ij}$ given by (2.2) are all equal to $c/n$ (for $n > c$). Thus any choice of vertex space gives the classical Erdős–Rényi random graph $G(n,c/n)$. (In fact, Erdős and Rényi [26, 27] studied the closely related graph $G(n,m)$ with a given number of edges, but as is well-known, the two random graph models are very similar (if we take $m = \lfloor cn/2 \rfloor$, say) and all our results on sizes of the largest components transfer easily from one model to the other.)

**Example 6.2** (the finite-type case)**.** Let $\mathcal{S} = \{s_1,\dots,s_r\}$ be finite. Then $\kappa$ is an $r \times r$ matrix. In this case, $G(n,\kappa)$ has vertices of $r$ different types (or colours), say $n_i$ vertices of type $i$, with two vertices of types $i$ and $j$ joined by an edge with probability $n^{-1}\kappa(i,j)$ (for $n \ge \max\kappa$). The condition (iv) means that $n_i/n \to \mu_i$ for each $i$ (in probability if the $n_i$ are random), where $\mu_i := \mu\{i\} \ge 0$.

This case has been studied by Söderberg [41; 42; 43; 44], who noted our Theorem 9.1 in this case (with $\kappa_n = \kappa$ for all $n$).

Note that Example 2.3 is equivalent to an example of this type.

**Remark 6.3.** Most of our proofs are based on a disguised form of the finite case, described by the following definition. (This includes Example 2.3 as a typical case.)

A kernel $\kappa$ on a (generalized) vertex space $(\mathcal{S},\mu,(\mathbf{x}_n))$ is *regular finitary* if $\mathcal{S}$ has a finite partition into $\mu$-continuity sets $S_1,\dots,S_r$ such that $\kappa$ is constant on each $S_i\times S_j$.

Clearly, if $\kappa$ is regular finitary on $(\mathcal{S},\mu,(\mathbf{x}_n))$ then the random graph $G^{\mathcal{V}}(n,\kappa)$ has the same distribution as a finite-type graph $G^{\mathcal{V}'}(n,\kappa')$, $\mathcal{V}'=(\mathcal{S}',\mu',(\mathbf{y}_n)_{n\ge1})$: take $\mathcal{S}'=\{1,\dots,r\}$, let $y_k=i$ whenever $x_k\in S_i$, and define $\mu'\{i\}$ and $\kappa'(i,j)$ in the obvious way.

A finite-type or regular finitary kernel $\kappa$ on a vertex space $\mathcal{V}$ is automatically graphical on $\mathcal{V}$.

**Example 6.4** (i.i.d. vertices)**.** For any separable metric space $\mathcal{S}$ equipped with a probability measure $\mu$, we can construct a vertex space by take $x_1,\dots,x_n$ to be i.i.d. random points in $\mathcal{S}$ with distribution $\mu$. (This has been proposed by, for example, Söderberg [41].) This is a standard setting, and various examples are obtained by choosing suitable kernels $\kappa$.

**Example 6.5** (Poisson process graph)**.** For any separable metric space $\mathcal{S}$ equipped with a probability measure $\mu$, and any $\lambda>0$, we can construct a generalized vertex space by letting $x_1,\dots,x_N$ be the points of a Poisson process on $\mathcal{S}$ with intensity measure $\lambda\mu$. In other words, $N$ has a Poisson distribution $\mathrm{Po}(\lambda)$, and, given $N$, the points are i.i.d. as in the preceding example. For any kernel $\kappa$ on $\mathcal{S}$, we then a random graph $\widetilde{G}_\lambda(\kappa)$, where we use $\lambda$ as a parameter instead of $n$.

Conditioned on $N=m$, this random graph is just $G(m,\tilde\kappa)$, with $x_1,\dots,x_m$ as in the preceding example and $\tilde\kappa:=(m/\lambda)\kappa$.

## 7. Graph limits

A very similar construction of random graphs has been used by Lovász and Szegedy [34] in their study of graph limits, see further e.g. Borgs, Chayes, Lovász, Sós and Vesztergombi [13] for further developments. More precisely, these authors assume that the kernel $\kappa$ satisfies $0\le\kappa\le1$ and consider the random graph defined as above, with $x_i$ i.i.d. on $\mathcal{S}$ as in Example 6.4, and usually taking $\mathcal{S}=[0,1]$ with Lebesgue measure, but they take

$$p_{ij}=\kappa(x_i,x_j) \tag{7.1}$$

instead of our (2.2). For each finite $n$, this gives the same random graphs as our definition (with this vertex space), by changing the notation and rescaling $\kappa$, but the behaviours as $n\to\infty$ are different; we consider the sparse case where the average degrees stay bounded and the number of edges grows like $n$, while they study the dense case where the number of edges grows like $n^2$.

This version of the construction, using (7.1) can also be used with an *infinite* number of vertices. By a theorem by Aldous [3] and Hoover [29], the resulting random graphs are exactly the exchangeable infinite random graphs. See further Diaconis and Janson [21] for the connection between exchangeable infinite random graphs and the theory of graph limits.

## 8. More definitions

We are going to state our main results in the following sections, but first we need some more definitions. For some of the results we need one additional condition.

**Definition.** A kernel $\kappa$ on a (generalized) vertex space $(\mathcal{S},\mu,(\mathbf{x}_n))$ is *reducible* if $\exists A\subset\mathcal{S}$ with $0<\mu(A)<\mu(\mathcal{S})$ such that $\kappa=0$ a.e. on $A\times(\mathcal{S}\setminus A)$; otherwise $\kappa$ is *irreducible*.

Thus $\kappa$ is irreducible if $A\subseteq\mathcal{S}$ with $\kappa=0$ a.e. on $A\times(\mathcal{S}\setminus A)$ implies that $\mu(A)=0$ or $\mu(\mathcal{S})$.

Roughly speaking, $\kappa$ is reducible if the vertex set of $G^{\mathcal{V}}(n,\kappa)$ can be split into two parts so that the probability of an edge from one part to the other is zero, and irreducible otherwise.

**A branching process.** A main tool to study components of $G(n,\kappa)$ is a branching process approximation, generalizing the classical branching process approximation for $G(n,p)$ (see, e.g., [31]). We use the multi-type Galton–Watson branching process with type space $\mathcal{S}$, where a particle of type $x\in\mathcal{S}$ is replaced in the next generation by a set of particles distributed as a Poisson process on $\mathcal{S}$ with intensity $\kappa(x,y)\,d\mu(y)$. (Thus, the number of children with types in a subset $A\subseteq\mathcal{S}$ has a Poisson distribution with mean $\int_A\kappa(x,y)\,d\mu(y)$, and these numbers are independent for disjoint sets $A$ and for different particles.) We denote this branching process, started with a single particle of type $x$, by $\mathfrak{X}_\kappa(x)$.

Let $\rho(\kappa;x)$ be the probability that this branching process $\mathfrak{X}_\kappa(x)$ survives for eternity.

We further define

$$\rho(\kappa):=\int_{\mathcal{S}}\rho(\kappa;x)\,d\mu(x).$$

When $\mu$ is a probability measure, this is the survival probability for the branching process above started with a single random point with distribution $\mu$.

**An integral operator.** Let $T_\kappa$ be the integral operator on $(\mathcal{S},\mu)$ with kernel $\kappa$, defined by

$$(T_\kappa f)(x)=\int_{\mathcal{S}}\kappa(x,y)f(y)\,d\mu(y),$$

for any (measurable) function $f$ such that this integral is defined (finite or $+\infty$) for a.e. $x$. Note that $T_\kappa f$ is defined for every $f\ge0$, with $0\le T_\kappa f\le\infty$.

If $\kappa \in L^1(\mathcal{S}\times\mathcal{S})$, as we assume, then $T_\kappa f$ is also defined for every bounded $f$; in this case $T_\kappa f \in L^1(\mathcal{S})$ and thus $T_\kappa f$ is finite a.e.

We define

$$\|T_\kappa\| := \sup\bigl\{\|T_\kappa f\|_2 : f \ge 0,\, \|f\|_2 \le 1\bigr\} \le \infty.$$

**Remark 8.1.** When finite, $\|T_\kappa\|$ is the norm of $T_\kappa$ as an operator in $L^2(\mathcal{S},\mu)$; it is infinite if $T_\kappa$ does not define a bounded operator in $L^2$.

Trivially, $\|T_\kappa\|$ is at most the Hilbert–Schmidt norm of $T_\kappa$:

$$\|T_\kappa\| \le \|T_\kappa\|_{HS} := \|\kappa\|_{L^2(\mathcal{S}\times\mathcal{S})} = \biggl(\iint_{\mathcal{S}^2} \kappa(x,y)^2\,d\mu(x)\,d\mu(y)\biggr)^{1/2}.$$

In particular, $\|T_\kappa\| < \infty$ if $\kappa$ is bounded.

**Remark 8.2.** Since $T_\kappa$ is symmetric, $\|T_\kappa\|$ equals also the *spectral radius* of the operator $T_\kappa$.

Define the non-linear operator $\Phi_\kappa$ by

$$\Phi_\kappa f := 1 - e^{-T_\kappa f}$$

for $f \ge 0$. For such $f$ we have $0 \le T_\kappa f \le \infty$, and thus $0 \le \Phi_\kappa f \le 1$. We can characterize $\rho(\kappa;x)$, and thus $\rho(\kappa)$, in terms of the non-linear operator $\Phi_\kappa$ as follows.

There is a (necessarily unique) maximum solution $\tilde\rho_\kappa$ to

$$\Phi_\kappa(\tilde\rho_\kappa) = \tilde\rho_\kappa, \tag{8.1}$$

i.e., a solution that pointwise dominates all other solutions. Furthermore, the survival probability $\rho(\kappa;x) = \tilde\rho_\kappa(x)$ for a.e. $x$, and

$$\Phi_\kappa(\rho_\kappa) = \rho_\kappa \quad a.e.,$$

where the function $\rho_\kappa$ is defined by $\rho_\kappa(x) := \rho(\kappa;x)$.

If $\|T_\kappa\| \le 1$, then $\tilde\rho_\kappa$ is identically zero, and this is thus the only solution. If $\|T_\kappa\| > 1$, then $\tilde\rho_\kappa$ is positive on a set of positive measure. Thus $\rho(\kappa) > 0$ if and only if $\|T_\kappa\| > 1$. In other words, the branching process $\mathfrak{X}_\kappa$ (with a random starting point) a.s. dies out if and only if $\|T_\kappa\| \le 1$.

Furthermore, If $\|T_\kappa\| > 1$ and $\kappa$ is irreducible, then $\tilde\rho_\kappa$ is the unique non-zero solution of (8.1), and $\tilde\rho_\kappa = \rho_\kappa > 0$ a.e.

**Example 8.3** (the Erdős-Rényi random graph)**.** In the Erdős–Rényi case, Example 6.1, the simplest choice is to let $\mathcal{S}$ consist of a single point. Then the operator $T_\kappa$ is simply multiplication by $c$, so $\|T_\kappa\| = c$ and Corollary 9.2 below yields the classical result that there is a phase transition at $c = 1$. Furthermore, the multi-type branching process $\mathfrak{X}_c$ below reduces in this case to a single-type Galton–Watson process with Poisson offspring distribution $\mathrm{Po}(c)$, and the function $\rho(c;x)$ below reduces to the single value $\rho(c)$, which is the survival probability $\rho(c)$ of this branching process and is given by the formula

$$\rho(c) = 1 - e^{-c\rho(c)}, \qquad \text{with } \rho(c) > 0 \text{ if } c > 1. \tag{8.2}$$

Consequently, in this case Theorem 9.1 on the size of the giant component reduces to the classical result of Erdős and Rényi [26]. (Note that (8.1) in this case reduces to (8.2).)

**Example 8.4** (the finite-type case)**.** In the finite-type case Example 6.2, the operator $T_\kappa$ is given by the matrix $(\kappa(i,j)\mu_j)_{ij}$, and $\|T_\kappa\|$ may be found by calculating the largest eigenvalue of this matrix.

## 9. Giant component

We denote the orders of the components of a graph $G$ by $C_1(G) \geq C_2(G) \geq \dots$, with $C_j(G) = 0$ if $G$ has fewer than $j$ components. We let $N_k(G)$ denote the total number of vertices in components of order $k$, and write $N_{\geq k}(G)$ for $\sum_{j\geq k} N_j(G)$, the number of vertices in components of order at least $k$. Our results are asymptotic, and all unspecified limits are taken as $n \to \infty$.

**Theorem 9.1.** *Let $(\kappa_n)$ be a graphical sequence of kernels on a (generalized) vertex space $\mathcal{V}$ with limit $\kappa$.*

(i) *If $\|T_\kappa\| \leq 1$, then $C_1\bigl(G^{\mathcal{V}}(n,\kappa_n)\bigr) = o_p(n)$, while if $\|T_\kappa\| > 1$, then $C_1\bigl(G^{\mathcal{V}}(n,\kappa_n)\bigr) = \Theta(n)$ whp.*

(ii) *For any $\varepsilon > 0$, whp we have*

$$\frac{1}{n} C_1\bigl(G^{\mathcal{V}}(n,\kappa_n)\bigr) \leq \rho(\kappa) + \varepsilon.$$

(iii) *If $\kappa$ is irreducible, then*

$$\frac{1}{n} C_1(G^{\mathcal{V}}(n,\kappa_n)) \xrightarrow{\mathrm{p}} \rho(\kappa).$$

*In all cases $\rho(\kappa) < 1$; furthermore, $\rho(\kappa) > 0$ if and only if $\|T_\kappa\| > 1$.*

As customary, we say that a sequence of random graphs $G_n$ (with $n$ vertices in $G_n$) has a *giant component* if $C_1(G_n) = \Theta(n)$ whp.

**Corollary 9.2.** *Let $\kappa$ be a graphical kernel on a (generalized) vertex space $\mathcal{V}$, and consider the random graphs $G^{\mathcal{V}}(n, c\kappa)$ where $c > 0$ is a constant. Then the threshold for the existence of a giant component is $c = \|T_\kappa\|^{-1}$. More precisely, if $c \leq \|T_\kappa\|^{-1}$, then $C_1\bigl(G^{\mathcal{V}}(n,c\kappa)\bigr) = o_p(n)$, while if $c > \|T_\kappa\|^{-1}$ and $\kappa$ is irreducible, then $C_1\bigl(G^{\mathcal{V}}(n,c\kappa)\bigr) = \rho(c\kappa)n + o_p(n) = \Theta_p(n)$.*

**Corollary 9.3.** *Let $\kappa$ be a graphical kernel on a (generalized) vertex space $\mathcal{V}$. Then the property that $G^{\mathcal{V}}(n,c\kappa)$ has whp a giant component holds for every $c > 0$ if and only if $\|T_\kappa\| = \infty$. Otherwise it has a finite threshold $c_0 > 0$.*

In the light of the results above, we say that a kernel $\kappa$ is *subcritical* if $\|T_\kappa\| < 1$, *critical* if $\|T_\kappa\| = 1$, and *supercritical* if $\|T_\kappa\| > 1$. We use the same expressions for a random graph $G(n,\kappa)$ and a branching process $\mathfrak{X}_\kappa$; this agrees with the standard notation for branching processes.

The number of edges in the graph at the point where the giant component emerges is maximal in the classical Erdős–Rényi case. (For normalization, we consider vertex spaces only.)

**Proposition 9.4.** *Let $\kappa_n$ be a graphical sequence of kernels on a vertex space $\mathcal{V}$ with limit $\kappa$, and assume that $\kappa$ is critical, i.e. $\|T_\kappa\| = 1$. Then $\frac{1}{n}e(G^{\mathcal{V}}(n,\kappa_n)) \xrightarrow{\mathrm{p}} \frac{1}{2}\iint \kappa \le 1/2$, with equality in the uniform case $\kappa = 1$.*

(This should not be surprising, since an inhomogeneity typically means that some part of the graph has higher density than the average, which makes it easier to creat a large component there.)

We can also determine the asymptotic number of edges in the giant component. As this is not always uniquely defined, for any graph $G$, let $\mathcal{C}_1(G)$ be the *largest component* of $G$, i.e., the component with most vertices, chosen according to any fixed rule if there is a tie. In order to state the next result concisely, let

$$\zeta(\kappa) := \frac{1}{2}\iint_{\mathcal{S}^2} \kappa(x,y)\bigl(\rho(\kappa;x) + \rho(\kappa;y) - \rho(\kappa;x)\rho(\kappa;y)\bigr)\,d\mu(x)\,d\mu(y).$$

**Theorem 9.5.** *Let $(\kappa_n)$ be a graphical sequence of kernels on a (generalized) vertex space $\mathcal{V}$ with irreducible limit $\kappa$. Then*

$$\frac{1}{n}e\bigl(\mathcal{C}_1(G^{\mathcal{V}}(n,\kappa_n))\bigr) \xrightarrow{\mathrm{p}} \zeta(\kappa).$$

The giant component is whp unique when it exists; the second largest component is much smaller. Indeed, only $o_p(n)$ vertices are in 'large' components other than the largest.

**Theorem 9.6.** *Let $(\kappa_n)$ be a graphical sequence of kernels on a (generalized) vertex space $\mathcal{V}$ with irreducible limit $\kappa$, and let $G_n = G^{\mathcal{V}}(n,\kappa_n)$. If $\omega(n) \to \infty$, then*

$$\sum_{j\ge 2:\ C_j(G_n)\ge\omega(n)} C_j(G_n) = o_p(n).$$

*In particular,*

$$C_2(G_n) = o_p(n).$$

**Remark 9.7.** If $\kappa$ and $\kappa'$ are two kernels with $\kappa' \ge \kappa$, then one can couple the corresponding graphs or branching processes so that $G(n,\kappa) \subseteq G(n,\kappa')$ or $\mathfrak{X}_\kappa \subseteq \mathfrak{X}_{\kappa'}$. Thus $\rho(\kappa) \le \rho(\kappa')$.

If $\kappa$ is irreducible and $\rho(\kappa) > 0$, then $\rho(\kappa') > \rho(\kappa)$ unless $\kappa' = \kappa$ a.e.

Similarly, the threshold $c_0(\kappa') := \|T_{\kappa'}\|^{-1}$ is at most $c_0(\kappa) := \|T_\kappa\|^{-1}$.

Here, however, somewhat surprisingly, we may have $c_0(\kappa') = c_0(\kappa)$ even if $\kappa' > \kappa$, see Section 18 for an example. On the other hand, it is easily seen that if $T_\kappa$ is compact and $\kappa' > \kappa$ on a set of positive measure, then $\|T_{\kappa'}\| > \|T_\kappa\|$ and thus $c_0(\kappa') < c_0(\kappa)$.

## 10. Stability

The giant component of $G_n = G^{\mathcal{V}}(n, \kappa_n)$ is stable in the sense that its size does not change much if we add or delete a few vertices or edges. Note that the vertices or edges added or deleted do not have to be random or independent of the existing graph; they can be chosen by an adversary after inspecting the whole of $G_n$.

**Theorem 10.1.** *Let $(\kappa_n)$ be a graphical sequence of kernels on a (generalized) vertex space $\mathcal{V}$ with irreducible limit $\kappa$, and let $G_n = G^{\mathcal{V}}(n, \kappa_n)$. For every $\varepsilon > 0$ there is a $\delta > 0$ (depending on $\kappa$) such that, whp,*

$$(\rho(\kappa) - \varepsilon)n \le C_1(G_n') \le (\rho(\kappa) + \varepsilon)n$$

*for every graph $G_n'$ that may be obtained from $G_n$ by deleting at most $\delta n$ vertices and their incident edges, and then adding or deleting at most $\delta n$ edges.*

In particular, if $G_n'$ is a graph on $[n] = 1, \dots, n$ with $e(G_n' \triangle G_n) = o_p(n)$ then

$$C_1(G_n') = C_1(G_n) + o_p(n) = \rho(\kappa)n + o_p(n).$$

As pointed out by Britton and Martin-Löf [15], the theorem has the following interpretation: suppose that $G_n$ represents the network of contacts that may allow the spread of an infectious disease from person to person, and that we wish to eliminate the possibility of an epidemic by vaccinating some of the population. Even if the entire network of contacts is known, if the source of the infection is not known, a significant (constant, as $n \to \infty$) proportion of the population must be vaccinated: otherwise, there is still a giant component in the graph on the unvaccinated people, and if the infection starts at one of its vertices, it spreads to $\Theta(n)$ people.

## 11. More on the phase transition

Fix a graphical kernel $\kappa$ on a (generalized) vertex space $\mathcal{V}$, and study $G^{\mathcal{V}}(n, c\kappa)$ for a real parameter $c > 0$ as in the corollary above.

By Theorem 9.1, the size of the largest component of $G^{\mathcal{V}}(n, c\kappa)$ is described by the function $\rho(c\kappa)$, which is 0 for $c \le c_0 := \|T_\kappa\|^{-1}$ and strictly positive for larger $c$. With $\mathcal{V}$ and $\kappa$ fixed, let us denote this function by $\rho(c)$, $c > 0$. It turns out that $\rho(c)$ is continuous on $(0, \infty)$.

Since $\rho(c) = 0$ for $c \le c_0$ but not for larger $c$, the function $\rho$ is not analytic at $c_0$; in physical terminology, *there is a phase transition at $c_0$*.

For the classical Erdős–Rényi random graph $G(n, c/n)$ (obtained with $\kappa = 1$), it is well-known that $\rho$ is continuous but the first derivative has a jump at $c_0 = 1$; more precisely, $\rho'$ jumps from 0 to $\rho'_+(c_0) = 2$. For finite $d$, we shall say that the phase transition in $G^{\mathcal{V}}(n, \kappa)$ has *exponent $d$* if $\rho(c_0 + \varepsilon) = \Theta(\varepsilon^d)$ as $\varepsilon \searrow 0$. As we have just noted, in $G(n, c/n)$ the phase transition has exponent 1. If $\rho(c_0 + \varepsilon) = o(\varepsilon^d)$ for all $d$, we say that

the phase transition has *infinite exponent*. We are deliberately avoiding the physical term 'order', as it is not used in a consistent way in this context.

It has been shown (Dorogovtsev, Mendes and Samukhin (2001), Durrett (2003) and Bollobás, Janson and Riordan (2005)) that in the case $\mathcal{S} = (0,1]$ and $\kappa(x,y) = 1/(x \vee y)$, the phase transition 'is of infinite order', i.e., has infinite exponent. We shall later see in that it is also possible to have a phase transition with any finite exponent larger than 1 (including non-integer values).

The next theorem shows that the phase transition has exponent 1 for a wide class of kernels $\kappa$, including all bounded $\kappa$.

**Theorem 11.1.** *Let $\kappa$ be a kernel on a (generalized) vertex space $(\mathcal{S}, \mu, (\mathbf{x}_n))$. Suppose that $\kappa$ is irreducible, and that*

$$\sup_x \int_{\mathcal{S}} \kappa(x,y)^2 \, d\mu(y) < \infty.$$

(i) *The function $c \mapsto \rho(c) := \rho(c\kappa)$ is analytic except at $c_0 := \|T_\kappa\|^{-1}$.*
(ii) *Furthermore, $T_\kappa$ has an eigenfunction $\psi$ of eigenvalue $\|T_\kappa\| < \infty$, and every such eigenfunction is bounded and satisfies*

$$\rho(c_0+\varepsilon) = 2c_0^{-1} \frac{\int_{\mathcal{S}} \psi \int_{\mathcal{S}} \psi^2}{\int_{\mathcal{S}} \psi^3} \varepsilon + O(\varepsilon^2), \qquad \varepsilon > 0, \tag{11.1}$$

*so $\rho'_+(c_0) = 2c_0^{-1} \int_{\mathcal{S}} \psi \int_{\mathcal{S}} \psi^2 / \int_{\mathcal{S}} \psi^3 > 0$ and $\rho$ has a phase transition at $c_0$ with exponent 1.*

**Corollary 11.2.** *Let $\kappa$ be an irreducible kernel such that (11.1) holds, and let $c_0 := \|T_\kappa\|^{-1} > 0$. Then $c_0 \rho'_+(c_0) \le 2$, with equality in the classical Erdős–Rényi case.*

## 12. Bounds on the small components

For the classical random graph $G(n, c/n)$ it is well-known that in the subcritical ($c < 1$) case, $C_1 = O(\log n)$ whp, and that in the supercritical ($c > 1$) case, $C_2 = O(\log n)$ whp. If we add some conditions, we obtain similar results. As before, we write $G_n$ for $G^{\mathcal{V}}(n, \kappa_n)$.

**Theorem 12.1.** *Let $(\kappa_n)$ be a graphical sequence of kernels on a (generalized) vertex space $\mathcal{V}$ with limit $\kappa$.*

(i) *If $\kappa$ is subcritical, i.e., $\|T_\kappa\| < 1$, and $\sup_{x,y,n} \kappa_n(x,y) < \infty$, then $C_1(G_n) = O(\log n)$ whp.*
(ii) *If $\kappa$ is supercritical, i.e., $\|T_\kappa\| > 1$, $\kappa$ is irreducible, and either $\inf_{x,y,n} \kappa_n(x,y) > 0$ or $\sup_{x,y,n} \kappa_n(x,y) < \infty$, then $C_2(G_n) = O(\log n)$ whp.*

Note that in part (ii) we draw the same conclusion from the very different assumptions $\inf_{x,y,n} \kappa_n(x,y) > 0$ and $\sup_{x,y,n} \kappa_n(x,y) < \infty$. There is no similar result for the subcritical case assuming only that $\inf_{x,y,n} \kappa_n(x,y) > 0$. Example: The random graph $G_n^{1/j}(c)$ defined in Example 2.2 with $0 < c <$

$1/4$ is subcritical and satisfies $C_1(G_n^{1/j}(c)) = n^{\Theta(1)}$ whp, see Bollobás, Janson and Riordan [5].

## 13. Degree sequence

The degree of a vertex of a given type $x$ is asymptotically Poisson with a mean

$$\lambda(x) := \int_{\mathcal{S}} \kappa(x,y)\, d\mu(y)$$

that depends on $x$. This leads to a mixed Poisson distribution for the degree $D$ of a random vertex of $G^{\mathcal{V}}(n,\kappa_n)$. We write $Z_k$ for the number of vertices of $G^{\mathcal{V}}(n,\kappa_n)$ with degree $k$.

**Theorem 13.1.** *Let $(\kappa_n)$ be a graphical sequence of kernels on a (generalized) vertex space $\mathcal{V}$ with limit $\kappa$. Define $\lambda(x)$ as above, and let $\Xi$ have the mixed Poisson distribution $\int_{\mathcal{S}} \mathrm{Po}(\lambda(x))\, d\mu(x)/\mu(\mathcal{S})$. Then, for any fixed $k \ge 0$,*

$$Z_k/n \xrightarrow{\mathrm{p}} \mathbb{P}(\Xi = k) = \int_{\mathcal{S}} \frac{\lambda(x)^k}{k!} e^{-\lambda(x)}\, d\mu(x).$$

In other words, if $D$ is the degree of a random vertex of $G^{\mathcal{V}}(n,\kappa_n)$, and we normalize so that $\mu(\mathcal{S}) = 1$, then

$$\mathcal{L}(D \mid G^{\mathcal{V}}(n,\kappa_n)) \xrightarrow{\mathrm{p}} \mathcal{L}(\Xi) = \int_{\mathcal{S}} \mathrm{Po}(\lambda(x))\, d\mu(x).$$

## 14. Distances between vertices

Let us write $d(v,w)$ for the graph distance between two vertices of $G_n = G^{\mathcal{V}}(n,\kappa_n)$, which we take to be infinite if they lie in different components.

We first consider the distance between two random vertices. In the subcritical and critical cases, when all components are small, this distance is whp $\infty$. In the supercritical case, we have the following result. Equivalently, the distance between two random vertices in the giant component is whp $(1+o(1))\log n/\log\|T_\kappa\|$.

**Theorem 14.1.** *Let $\kappa_n$ be a graphical sequence of kernels on a (generalized) vertex space $\mathcal{V}$ with limit $\kappa$, with $\|T_\kappa\| > 1$. Let $G_n = G^{\mathcal{V}}(n,\kappa_n)$, and let $v$ and $w$ be two independently chosen random vertices in $G_n$.*

(i) *If $\kappa$ is irreducible and $1 < \|T_\kappa\| < \infty$, then*

$$\bigl(d(v,w)/\log n \mid d(v,w) < \infty\bigr) \xrightarrow{\mathrm{p}} 1/\log\|T_\kappa\|$$

(ii) *If $\kappa$ is irreducible and $\|T_\kappa\| = \infty$, then there is a function $f(n) = o(\log n)$ such that*

$$\mathbb{P}\bigl(d(v,w) \le f(n) \mid d(v,w) < \infty\bigr) \to 1.$$

Secondly, we consider the maximal finite distance. We define the 'diameter' of $G_n$ as

$$\operatorname{diam}(G_n) := \max\{d(v,w) \,:\, v,w \in V(G),\, d(v,w) < \infty\},$$

the maximum of the diameters of the components of $G_n$.

**Theorem 14.2.** *Let $\kappa$ be a kernel on a finite (generalized) vertex space $(\mathcal{S},\mu,(\mathbf{x}_n)_{n\ge1})$, $\mathcal{S} = \{1,2,\dots,r\}$, with $\mu(\{i\}) > 0$ for each $i$. If $0 < \|T_\kappa\| < 1$, then*

$$\frac{\operatorname{diam}(G_n)}{\log n} \xrightarrow{\mathrm{p}} \frac{1}{\log \|T_\kappa\|^{-1}}$$

*as $n \to \infty$, where $G_n = G^{\mathcal{V}}(n,\kappa)$. If $\|T_\kappa\| > 1$ and $\kappa$ is irreducible, then*

$$\frac{\operatorname{diam}(G_n)}{\log n} \xrightarrow{\mathrm{p}} \frac{2}{\log \|T_{\hat\kappa}\|^{-1}} + \frac{1}{\log \|T_\kappa\|},$$

*where $\hat\kappa$ is the dual kernel to $\kappa$, defined by $\hat\kappa(x,y) = \kappa(x,y)$ on $(\mathcal{S},\hat\mu)$, where the measure $\hat\mu$ is defined by*

$$d\hat\mu(x) = (1 - \rho(\kappa;x))\, d\mu(x).$$

Thus, in the supercritical case, the maximal finite distance is of the same order $\log n$ as the typical finite distances, but the constants differ.

## 15. Even more examples

**Example 15.1** (The homogeneous case)**.** More generally, let the vertex space $(\mathcal{S},\mu,(\mathbf{x}_n)_{n\ge1})$ be arbitrary, and let $\kappa$ be irreducible and such that $\int_{\mathcal{S}} \kappa(x,y)\, d\mu(y)$ is independent of $x \in \mathcal{S}$, i.e., that

$$\int_{\mathcal{S}} \kappa(x,y)\, d\mu(y) = c \qquad \text{for every } x,$$

for some constant $c$. (This says roughly that, asymptotically, all vertices have the same average degree.)

Then $T_\kappa 1 = c$, so the constant function 1 is a positive eigenfunction with eigenvalue $c$, and thus $\|T_\kappa\| = c$, and by Theorem 9.1 there is a giant component (and $\rho(\kappa) > 0$) if and only if $c > 1$.

In the branching process $\mathfrak{X}_\kappa$, the number of children of each particle has a $\mathrm{Po}(c)$ distribution. Hence, ignoring the types of the particles, the distributions of the process $\mathfrak{X}_\kappa$ and the single-type process with $\mathrm{Po}(c)$ offspring are the same. In particular, $\rho(\kappa) = \rho(c)$, and $\rho(\kappa) = \rho(c)$ is given by the same equation (8.2) as in the Erdős–Rényi case Example 6.1.

Thus, the global behaviour of $G(n,\kappa)$ is exactly the same as that of $G(n,c/n)$, at least in terms of the size of the giant component. The local behaviour can be quite different, though. For example, $G(n,\kappa)$ may have many more triangles or other small cycles than $G(n,c/n)$. On the other hand, the vertex degrees have an asymptotic $\mathrm{Po}(c)$ distribution just as in $G(n,c/n)$.

A natural example of such a homogeneous $\kappa$ is given by taking $\mathcal{S}$ as $(0,1]$ (now better regarded as the circle $\mathbb{T}$), $\mu$ as Lebesgue measure, and $\kappa(x,y) = h(x-y)$ for an even function $h \ge 0$ of period 1. For example, $h$ can be constant on a small interval $(-\delta,\delta)$ and vanish outside it; this gives a modification of $G(n,c/n)$ where only "short" edges are allowed.

More generally, $\mathcal{S}$ can be any compact homogeneous space, for example a sphere, with Haar measure $\mu$ and an invariant metric $d$, and $\kappa(x,y)$ a function of the distance $d(x,y)$.

**Example 15.2.** Take $\mathcal{S} = (0,1]$ with $\mu$ the Lebesgue measure, and let $x_i = i/n$. Set $\kappa(x,y) = \mathbf{1}[x+y \le 1]$ and consider the kernel $c\kappa$, so that

$$p_{ij} = \begin{cases} c/n, & i+j \le n; \\ 0, & i+j > n. \end{cases}$$

Thus $G(n,c\kappa)$ can be obtained from the random graph $G(n,c/n)$ by deleting all edges $ij$ with $i+j>n$.

The operator $T_\kappa$ is compact, and it easy to see that it has eigenvalues $(-1)^k\omega_k^{-1}$ and eigenfunctions $\cos(\omega_k x)$, with $\omega_k = (k+1/2)\pi$, $k = 0,1,\dots$. Hence $\|T_\kappa\| = 2/\pi$ and the critical value is $c_0 = \pi/2$. Theorem 11.1 shows that at the critical value we have $c_0\rho'_+(c_0) = 3/2$.

**Example 15.3** (Edge percolation)**.** Let $\kappa$ be an irreducible graphical kernel on a (generalized) vertex space $\mathcal{V}$ with $\|T_\kappa\| > 1$, and let $0 < p \le 1$. Independently of everything else, keep each edge in $G(n,\kappa)$ with probability $p$ and delete it with probability $1-p$. Denote the resulting graph by $G^{\langle p\rangle}(n,\kappa)$.

This random graph $G^{\langle p\rangle}(n,\kappa)$ is nothing but $G(n,\tilde\kappa_n)$, where

$$\tilde\kappa_n(x,y) := p\big(\kappa(x,y)\wedge n\big).$$

Clearly, $x_n \to x$ and $y_n \to y$ imply $\tilde\kappa_n(x_n,y_n) \to p\kappa(x,y)$, provided $(x,y)$ is a point of continuity of $\kappa$. Furthermore, $\mathbb{E}\, e\big(G^{\langle p\rangle}(n,\kappa)\big) = p\,\mathbb{E}\, e\big(G(n,\kappa)\big) \to p\frac12\iint\kappa$. Hence, Theorem 9.1 applies with $\kappa$ replaced by $p\kappa$, so

$$n^{-1}C_1\big(G^{\langle p\rangle}(n,\kappa)\big) \xrightarrow{\mathrm{p}} \rho(p\kappa).$$

In particular, $G^{\langle p\rangle}(n,\kappa)$ has whp a component of order $\Theta(n)$ if and only if $\|T_{p\kappa}\| > 1$, i.e., if $p > \|T_\kappa\|^{-1}$.

**Example 15.4** (Vertex percolation)**.** Independently of everything else, keep each vertex in $G(n,\kappa)$ with probability $p$ and delete it with probability $1-p$. Denote the resulting graph by $G^{[p]}(n,\kappa)$.

$G^{[p]}(n,\kappa)$ is the graph $G(m,\kappa_n)$ obtained from a random sample $\widetilde{x}_1,\dots,\widetilde{x}_m$ of the points $x_1,\dots,x_n$, rather than from all of them.

Theorem 9.1 applies, with $\mu$ replaced by $p\mu$, and it follows that

$$n^{-1}C_1\big(G^{[p]}(n,\kappa)\big) \xrightarrow{\mathrm{p}} p\rho(p\kappa).$$

In particular, $G^{[p]}(n,\kappa)$ has whp a component of order $\Theta(n)$ if and only if $\|T_{p\kappa}\|>1$, i.e. if $p>\|T_\kappa\|^{-1}$. We thus obtain the same threshold for vertex percolation in $G(n,\kappa)$ as for edge percolation.

In the final sections we study further examples, yielding random graph models that have been previously studied by various authors.

## 16. Dubins' model

A common setting is the following, encountered already in Example 2.1: the vertex space $\mathcal{V}$ is $(\mathcal{S},\mu,(\mathbf{x}_n)_{n\ge1})$, where $\mathcal{S}=(0,1]$, $\mu$ is the Lebesgue measure, and $\mathbf{x}_n=(x_1,\dots,x_n)$ with $x_i=i/n$. In this case, we have $p_{ij}=\kappa(i/n,j/n)/n\wedge1$ for the probability of an edge between vertices $i$ and $j$. We shall in this and the following sections consider several choices of $\kappa$ in some detail.

Observe first that if $\kappa$ is a positive function on $(0,\infty)^2$ that is homogeneous of degree $-1$, then (2.2) yields $p_{ij}=\kappa(i,j)\wedge1$. Since this does not depend on $n$, in this case we can also consider the infinite graph $G(\infty,\kappa)$, defined in the same way as $G_n=G^{\mathcal{V}}(n,\kappa)$ but on the vertex set $\{1,2,\dots\}$. Note that the graphs $G^{\mathcal{V}}(n,\kappa)$ are induced subgraphs of $G(\infty,\kappa)$ and that we can construct them by successively adding new vertices, and for each new vertex an appropriate random set of edges to earlier vertices.

We first consider, as in Example 2.2, $\kappa(x,y)=c/(x\vee y)$ with $c>0$, so that if $j\ge c$ then

$$p_{ij}=c/j \qquad \text{for } i<j. \tag{16.1}$$

In this case we can regard $G^{\mathcal{V}}(n,\kappa)$ as a sequence of graphs grown by adding new vertices one at a time where, when vertex $k$ is added, it gets $\mathrm{Bi}(k-1,c/k)$ edges, whose other endpoints are chosen uniformly among the other vertices. (We might instead take $\mathrm{Po}(c)\wedge(k-1)$ new edges, without any difference in the asymptotic results below.)

This infinite graph $G(\infty,\kappa)$ was considered by Dubins in 1984, who asked when $G(\infty,\kappa)$ is a.s. connected. Dubins' question was answered partially by Kalikow and Weiss [32]. A little later Shepp [40] proved that $G(\infty,\kappa)$ is a.s. connected if and only if $c>1/4$. This result was generalized to more general homogeneous kernels by Durrett and Kesten [25].

The finite random graph $G^{\mathcal{V}}(n,\kappa)$ with this $\kappa$, i.e., with edge probabilities given by (16.1), has been studied by Durrett [24], who points out that it has the same critical value $c=1/4$ for the emergence of a giant component as the infinite version has for connectedness, and by Bollobás, Janson and Riordan [5] who rigorously show that this example has a phase transition with infinite exponent. More precisely, denoting $\rho(\kappa)$ by $\rho(c)$, it was shown by Riordan [39] that

$$\rho(1/4+\varepsilon)=\exp\bigl(-\tfrac{\pi}{2}\varepsilon^{-1/2}+O(\log\varepsilon)\bigr). \tag{16.2}$$

A similar formula for the closely related CHKNS model (see Section 18), introduced by Callaway, Hopcroft, Kleinberg, Newman and Strogatz [17],

had been given earlier by Dorogovtsev, Mendes and Samukhin [23] using non-rigorous methods.

To find the critical value by our methods, we have to find the norm of $T_\kappa$ on $L^2(0,1)$. Using the isometry $U : f \mapsto e^{-x/2} f(e^{-x})$ of $L^2(0,1)$ onto $L^2(0,\infty)$, we may instead consider $\widetilde{T}_\kappa := UT_\kappa U^{-1}$, which by a simple calculation is the integral operator on $L^2(0,\infty)$ with kernel

$$\tilde\kappa(x,y) = e^{-x/2}\kappa(e^{-x},e^{-y})e^{-y/2} = ce^{-x/2-y/2+x\wedge y} = ce^{-|x-y|/2}.$$

Hence $\widetilde{T}_\kappa$ is the restriction to $(0,\infty)$ of the convolution with $h(x) := ce^{-|x|/2}$. Because of translation invariance, it is easily seen that $\widetilde{T}_\kappa$ has the same norm as convolution with $h$ on $L^2(-\infty,\infty)$, and taking the Fourier transform we find

$$\|T_\kappa\| = \|\widetilde{T}_\kappa\| = \|f \mapsto h * f\|_{L^2(-\infty,\infty)} = \sup_{\xi\in\mathbb{R}} |\hat h(\xi)| = \int_{-\infty}^{\infty} h(x)\,dx = 4c.$$

Thus, Theorem 9.1 shows that there is a giant component if and only if $c > 1/4$, as shown in Durrett [24] and [5].

To find the size of the giant component is more challenging, and we refer to Riordan [39] for a proof of (16.2). Note that Theorem 11.1(ii) does *not* apply. Indeed, it is easy to see that $T_\kappa$ is a non-compact operator, and that it has no eigenfunctions at all in $L^2$. We suspect that this is connected to the fact that the phase transition has infinite exponent.

## 17. The mean-field scale-free model

Another interesting case with a homogeneous kernel as in Section 16 is $\kappa(x,y) = c/\sqrt{xy}$ with $c > 0$; then, for $ij \ge c^2$, we have

$$p_{ij} = c/\sqrt{ij}. \tag{17.1}$$

This model has been studied in detail by Riordan [39]. Considering the sequence $G^{\mathcal{V}}(n,\kappa)$ as a growing graph, in this case, together with each new vertex we add a number of edges that has approximately a Poisson Po($2c$) distribution; the other endpoint of each edge is chosen with probability proportional to $i^{-1/2}$, which is approximately proportional to the degree of vertex $i$. Hence, this random graph model resembles the growth with preferential attachment model of Barabási and Albert [4], which was made precise as the *LCD model* by Bollobás and Riordan [10]; see also [39].

In fact, up to a factor of $1 + o(i^{-1})$ in the edge probabilities, the model defined by (17.1) is the so called 'mean-field' version of the Barabási–Albert model, having the same individual edge probabilities, but with edges present independently. (This by now common use of 'mean-field' is not the standard one in physics, where it normally means that all vertices interact equally.)

In this case, $T_\kappa$ is an unbounded operator, because $x^{-1/2} \notin L^2(0,1)$, and thus there is no threshold. In other words, $\rho(c) := \rho(\kappa) > 0$ for every $c > 0$.

As shown by Riordan [39], $\rho(c)$ grows very slowly at first in this case too; more precisely,

$$\rho(c) \sim 2e^{1-\gamma}\exp\bigl(-1/(2c)\bigr) \qquad \text{as } c \to 0, \tag{17.2}$$

where $\gamma$ is Euler's constant. The result in [39] for the Barabási–Albert model is different, showing that in this model the dependence between edges is important.

**Remark 17.1.** Random graphs related to the ones defined here and in Section 16 but with some dependence between edges (and thus not covered by the present paper) can be obtained by adding at each new vertex a number of edges with some other distribution, for example $\mathrm{Bi}(m,p)$ for some fixed $m$ and $p$. Such random graphs have been considered in [9; 11; 20; 39], and these papers show that not only the expected numbers of edges added at each step are important, but also the variances; the edge dependencies shift the threshold.

## 18. The CHKNS model

The CHKNS model of Callaway, Hopcroft, Kleinberg, Newman and Strogatz [17] grows from a single vertex; vertices are added one by one, and after each vertex is added, an edge is added with probability $\delta$; the endpoints are chosen uniformly among all existing vertices. (Multiple edges are allowed; this does not matter for the asymptotics.)

Following Durrett [24], we consider a modification (which is perhaps at least as natural): after adding each vertex, add a Poisson $\mathrm{Po}(\delta)$ number of edges to the graph, again choosing the endpoints of these edges uniformly at random. Thus, when vertex $k$ is added, each existing pair of vertices acquires $\mathrm{Po}\bigl(\delta/\binom{k}{2}\bigr)$ new edges, and these numbers are independent. When we have reached $n$ vertices, the number of edges between vertices $i$ and $j$, with $1 \le i \le j \le n$, is thus Poisson with mean

$$e_{ij} := \sum_{k=j}^{n} \frac{\delta}{\binom{k}{2}} = 2\delta \sum_{k=j}^{n} \frac{1}{k(k-1)} = 2\delta\Bigl(\frac{1}{j-1} - \frac{1}{n}\Bigr),$$

and the probability that there is one or more edges between $i$ and $j$ is $p_{ij} := 1 - \exp(-e_{ij})$.

Hence, ignoring multiple edges, we have a graph $G_n$ of our type, with $\mathcal{S} = (0,1]$, $\mu$ Lebesgue measure, $x_i = i/n$ and

$$\begin{aligned}\kappa_n(x,y) &:= n\Bigl(1 - \exp\Bigl(-2\delta\Bigl(\frac{1}{n(x\vee y)-1} - \frac{1}{n}\Bigr)\Bigr)\Bigr)\\ &\to \kappa(x,y) := 2\delta\Bigl(\frac{1}{x\vee y} - 1\Bigr).\end{aligned} \tag{18.1}$$

Theorem 9.1 shows that $C_1(G_n)/n \xrightarrow{\mathrm{p}} \rho(\kappa)$.

The original CHKNS model, $\widetilde{G}_n$, say, can be treated by a comparison argument. It follows that $C_1(\widetilde{G}_n)/n \xrightarrow{\mathrm{p}} \rho(\kappa)$ holds for the CHKNS model too.

In particular, the threshold for the CHKNS model, as well as for Durrett's modification, is given by $\|T_\kappa\| = 1$, or $2\delta = \|T\|^{-1}$, where $T$ is the integral operator with kernel $1/(x \vee y) - 1$ on $L^2(0,1)$. This kernel is strictly smaller that the kernel $1/(x \vee y)$ considered in Section 16. However, changing variables as there, we see that $T$ is equivalent to the operator on $L^2(0,\infty)$ with kernel $e^{-|x-y|/2} - e^{-(x+y)/2}$. Using translational invariance of the operator with kernel $e^{-|x-y|/2}$ considered in Section 16, it is easily seen that $T$ has the same norm as this operator, namely 4.

Thus the thresholds for the CHKNS model and Durrett's modification are both given by $2\delta = 1/4$, i.e. $\delta = 1/8$, as was found by non-rigorous arguments by Callaway, Hopcroft, Kleinberg, Newman and Strogatz [17] and Dorogovtsev, Mendes and Samukhin [23], and first proved rigorously by Durrett [24].

To study the size of the giant component in these models, let us write $\kappa_0(x,y) := 1/(x \vee y)$ and $\kappa_1(x,y) := 1/(x \vee y) - 1$. Then $\kappa_1 < \kappa_0$, and thus $\rho(c\kappa_1) \le \rho(c\kappa_0)$ for each $c > 0$; in fact, it is easy to see that we have strict inequality for $c > 1/4$, see Remark 9.7. Note, however, that, as pointed out by Durrett [24], we have the same threshold $1/4$ for both kernels although we have twice as many edges in $G(n, c\kappa_0)$ as in $G(n, c\kappa_1)$. On the other hand, it is easy to see by a coupling argument that for every $\eta$ with $0 < \eta < 1$,

$$\rho(c\kappa_0) \ge \rho(c\kappa_1) \ge \eta\rho\big((1-\eta)c\kappa_0\big).$$

Taking $c = 1/4 + \varepsilon$ and $\eta = \varepsilon^2$, relation (16.2) for $\rho(c\kappa_0)$ implies the same estimate for $\rho(c\kappa_1)$; in other words, if $\delta = 1/8 + \varepsilon$, then the size of the giant component is given by

$$\rho(\kappa) = \rho(2\delta\kappa_1) = \exp\Big(-\frac{\pi}{2\sqrt{2}}\varepsilon^{-1/2} + O(\log\varepsilon)\Big).$$

In particular, the phase transition has infinite exponent in this example too.

As noted in Section 16, a similar but more precise formula (with no error term, and a particular constant in front of the exponential) has been given by Dorogovtsev, Mendes and Samukhin [23] using non-rigorous methods; see also Durrett [24].

## 19. The rank 1 case

The *rank 1 case* is when the kernel $\kappa$ has the form $\kappa(x,y) = \psi(x)\psi(y)$ for some function $\psi > 0$ on $\mathcal{S}$. We shall assume that $\int \psi \, d\mu < \infty$, but not necessarily that $\int \psi^2 \, d\mu < \infty$; we further assume that the kernel is graphical.

This is a special case of our general model that, while very restrictive, is also very natural, and includes or is closely related to many random graph models considered by other authors.

The function $\psi(x)$ can be interpreted as the "activity" of a vertex at $x$, with the probability of an edge between two vertices proportional to the product of their activities.

In the rank 1 case, $T_\kappa f = \bigl(\int f\psi\bigr)\psi$, and $\|T_\kappa\| = \|\psi\|_2^2 = \int \psi^2\,d\mu \le \infty$. Thus $T_\kappa$ is bounded if and only if $\psi \in L^2$, in which case $T_\kappa$ has rank 1, so it is compact, and $\psi$ is the unique (up to multiplication by constants) eigenfunction with non-zero eigenvalue.

By Theorem 13.1, the distribution of vertex degrees is governed by the distribution of the function $\lambda(x) = (\int \psi\,d\mu)\psi(x)$ on $(\mathcal{S},\mu)$. In particular, the degree sequence will (asymptotically) have a power-law tail if the distribution of $\lambda(x)$ has; for example, if $\mathcal{S} = (0,1]$ with $\mu$ Lebesgue measure, and $\psi(x) = cx^{-1/p}$.

Another, perhaps more canonical, version is to take $\psi(x) = x$ on $\mathcal{S} = [0,\infty)$, with a suitable finite Borel measure $\mu$. Note that every random graph considered in this example may be defined in this way, since we may map $\mathcal{S}$ to $[0,\infty)$ by $x \mapsto \psi(x)$. Alternatively, we may map by $x \mapsto \lambda(x)$ and have $\psi(x) = cx$ with $c > 0$ and $\lambda(x) = x$.

Random graphs of this type have been studied by several authors, e.g, Chung and Lu [18], Norros and Reittu [38], and Britton, Deijfen and Martin-Löf [14]. Actually, in [18] and [38] the edge probabilities $p_{ij}$ are given by $p_{ij} := w_iw_j/\sum_{i=1}^n w_i$, with $w_i$ deterministic in [18] and random in [38]. Under suitable conditions on the $w_i$, these examples are also special cases of our general model. More precisely, if $w_1,\dots,w_n$ are deterministic or random such that their empirical distribution $\frac{1}{n}\sum_{i=1}^n \delta_{w_i} \xrightarrow{\mathrm{p}} \mu$ for some probability measure $\mu$ on $\mathcal{S} = [0,\infty)$, and further $\frac{1}{n}\sum_{i=1}^n w_i \xrightarrow{\mathrm{p}} \omega := \int_0^\infty x\,d\mu(x) \in (0,\infty)$, then we can take $x_i = w_i\bigl(\sum_j w_j/n\omega\bigr)^{-1/2} \in \mathcal{S}$ and $\psi(x) = \omega^{-1/2}x$, and thus $\kappa(x,y) = xy/\omega$. (These assumptions on $w_i$ are satisfied for example for random i.i.d. $w_i$, as in [38], if we further assume $\mathbb{E}\,w_i = \omega < \infty$, taking $\mu = \mathcal{L}(w_1)$.) It is easily seen that all conditions in Section 4 are satisfied, except possibly (vii). Typically, (vii) too is satisfied, but sometimes we may have to condition on e.g. $\sum_1^n w_i > \frac{1}{2}n\omega$ first, cf. Remark 4.2.

**Remark 19.1.** The random graphs $G(n,\kappa)$ obtained from rank 1 kernels should be compared to the random graphs with a given (suitably chosen) degree sequence $(d_i)_1^n$, studied by, for example, Luczak [35], Molloy and Reed [36; 37] and (in the power-law case) Aiello, Chung and Lu [1]. Note that in this model, the probability of an edge between $i$ and $j$ is roughly $d_id_j/n$, but there are dependencies between the edges. The results by Molloy and Reed [36; 37] on existence and size of a giant component in this model fits well with our results in the rank 1 case, although we see no strict implication.

To study the phase transition in the rank 1 case, let us now consider the kernel $c\kappa(x,y) = c\psi(x)\psi(y)$, with $c > 0$ a parameter. We study the size of the giant component (if any) as a function of $c$, and let $\alpha(c) := c\int \psi\rho_{c\kappa}\,d\mu$, where, as before, $\rho_{c\kappa}(x) = \rho(c\kappa;x)$ is the survival probability of

the branching process $\mathfrak{X}_{c\kappa}(x)$. Then $\rho_{c\kappa} \searrow 0$ a.e. as $c \searrow c_0$, and thus so, by dominated convergence,

$$\alpha(c)/c \searrow 0 \text{ as } c \searrow c_0.$$

We have $T_{c\kappa}\rho_{c\kappa} = cT_\kappa\rho_{c\kappa} = \alpha(c)\psi$, and thus

$$\rho_{c\kappa} = \Phi_{c\kappa}(\rho_{c\kappa}) = 1 - e^{-T_{c\kappa}\rho_{c\kappa}} = 1 - e^{-\alpha(c)\psi}.$$

Let

$$\beta(t) := \int_{\mathcal{S}} \Bigl(1 - e^{-t\psi(x)}\Bigr)\psi(x)\,d\mu(x), \qquad t \ge 0.$$

Then,

$$\alpha(c) = c\int_{\mathcal{S}} \rho_{c\kappa}\psi\,d\mu = c\beta\bigl(\alpha(c)\bigr).$$

so $c = \alpha(c)/\beta\bigl(\alpha(c)\bigr)$, i.e., $\alpha$ is the inverse function to $t \mapsto \gamma(t) := t/\beta(t)$.

Let us consider some concrete examples. Take, again, $\mathcal{S} = (0,1]$ with $\mu$ Lebesgue measure, and let $\psi(x) = x^{-1/p}$ where $1 < p \le \infty$. We shall use $C$, $C_1$, etc. to denote various positive constants that depend on $p$.

*Case 1:* $1 < p < 2$. In this case, $\|\psi\|_2 = \infty$, so $c_0 = 0$. Calculations yield

$$\rho(c) \sim C_3\alpha(c) = C_3\gamma^{-1}(c) \sim C_4 c^{1/(2-p)} \quad \text{as } c \to 0.$$

Note that this exponent $1/(2-p)$ may be any real number in $(1,\infty)$.

*Case 2:* $p = 2$. This is the case (17.1) studied in Section 17 and [39]. We still have $\|\psi\|_2 = \infty$ and thus $c_0 = 0$. We now find that

$$\rho(c) = e^{-(1+o(1))/2c} \qquad \text{as } c \to 0.$$

More refined estimates can be obtained in the same way, see (17.2) and [39].

*Case 3:* $2 < p < 3$. For $p > 2$ we have $\int \psi^2\,d\mu < \infty$, and thus $c_0 > 0$, so we have a phase transition. (In fact, $c_0 = 1 - 2/p$.) Calculations yield

$$\rho(c_0 + \varepsilon) \sim C_4\alpha(c_0 + \varepsilon) \sim C_5\varepsilon^{1/(p-2)} \qquad \text{as } \varepsilon \searrow 0.$$

We thus have a phase transition at $c_0$ with exponent $1/(p-2)$. Note that this exponent may be any real number in $(1,\infty)$. (Taking instead e.g. $\psi(x) = x^{-1/2}\ln^{-1}(e^3/x)$, it is similarly seen that there is a phase transition with infinite exponent.)

*Case 4:* $p = 3$. Similarly, with $c_0 = 1/3$,

$$\rho(c_0 + \varepsilon) \sim C\alpha(c_0 + \varepsilon) \sim C_1\varepsilon/\ln(1/\varepsilon) \qquad \text{as } \varepsilon \searrow 0,$$

so $\rho'(c_0) = 0$.

*Case 5:* $3 < p \le \infty$. In this case, $\int \psi^3\,d\mu < \infty$. We find

$$\rho(c_0 + \varepsilon) \sim C_2\alpha(c_0 + \varepsilon) \sim C_3\varepsilon,$$

so we have a phase transition with exponent 1. (This is similar to Theorem 11.1, although the conditions are not quite satisfied.)

**Remark 19.2.** In these examples with $\psi(x) = x^{-1/p}$, the degree distribution has by Theorem 13.1 a power-law tail with $\mathbb{P}(D > t) \sim t^{-p}$ as $t \to \infty$. In these examples, there is thus a connection between the exponent of the phase transition and the exponent of the tail of the degree distribution. It is tempting to guess that this connection, at least in some form, holds more generally. In this direction, note that for the rank 1 case, by Theorem 13.1, $\mathbb{P}(D > t) \sim ct^{-p}$ if and only if $\mu\{x : \psi(x) > t\} \sim c_1 t^{-p}$, and thus by Corollary 9.3, the threshold $c_0 = 0$ for $p \le 2$ while $c_0 > 0$ for $p > 2$.

On the other hand, in the case $p > 2$ where the asymptotic degree distribution has finite variance and Corollary 9.3 yields a positive threshold for the existence of a giant component, there is no reason to expect a more precise connection between the *existence* of a giant component for a given $\psi$ or $\kappa$ and the tail of the degree distribution. For example, we can create or destroy a giant component by multiplying $\psi$ by a suitable positive constant. Moreover, even if, say, $\lambda(x) \sim x^{-1/p}$ (without constant factor), this only gives the asymptotics of the distribution of $\psi$ on $(\mathcal{S}, \mu)$, while the existence of a giant component depends on whether $\|T_\kappa\| = \int \psi^2 \, d\mu > 1$ or not, which is determined by the bulk of the distribution of $\psi$ and is essentially independent of its tail behaviour. The same applies to the examples with given degree sequences in e.g. [1]; the relation between the tail of the degree distribution and existence of a giant component seen in these examples thus depends on the details of the particular distributions chosen, and does not reflect a general phenomenon.

## 20. Turova's model

Turova [45; 46; 47; 48] has studied a dynamical random graph $G(t)$, $t \ge 0$, defined as follows, using three parameters $\gamma > 0$, $\lambda > 0$ and $\delta \ge 0$. The graph starts with a single vertex at time $t = 0$. Each existing vertex produces new, initially isolated, vertices according to a Poisson process with intensity $\gamma$. As soon as there are at least two vertices, each vertex sends out edges according to another Poisson process with intensity $\lambda$; the other endpoint is chosen uniformly among all other existing vertices. (Multiple edges are allowed, but this makes little difference.) Vertices live for ever, but edges die with intensity $\delta$, i.e., the lifetime of an edge has an exponential distribution with mean $1/\delta$. (All these random processes and variables are independent.)

By homogeneity we may assume $\gamma = 1$; the general case follows by replacing $\lambda$ and $\delta$ by $\lambda/\gamma$ and $\delta/\gamma$ and changing the time scale.

The vertices proliferate according to a Yule process (binary fission process): writing $N(t)$ for the number of vertices at time $t$, the probability that a new vertex is added in the infinitesimal time interval $[t, t + dt]$ is $N(t)\,dt$. It is well-known that

$$e^{-t} N(t) \xrightarrow{\text{a.s.}} W \qquad \text{as } t \to \infty \tag{20.1}$$

for a random variable $W$ with $W > 0$ a.s. (In fact, $W \sim \mathrm{Exp}(1)$.)

We condition on the vertex process, and assume, as we may by (20.1), that

$$e^{-t}N(t) \to w \qquad \text{as } t \to \infty \tag{20.2}$$

for some $w > 0$. We then take $\mathcal{S} = [0,\infty)$ and let $x_1, \dots, x_{N(t)}$ be the ages of the particles existing at time $t$. It is easily checked that this gives a vertex space $(\mathcal{S}, \mu, \mathbf{x}_n)$, where $\mu$ is the measure on $[0,\infty)$ given by $d\mu/dx = e^{-x}$ (the exponential distribution).

Moreover, the number of edges at time $t$ between two vertices of ages $x_i$ and $x_j$ has a Poisson distribution with mean $\kappa_t^*(x,y)/N(t)$, where

$$\kappa_t^*(x,y) := 2\lambda \int_0^{x\wedge y} e^{-\delta s} \frac{N(t)}{N(t-s)-1}\, ds.$$

It is easily checked that if $\delta \neq 1$ and $x_t \to x$, $y_t \to y$, then

$$\kappa_t^*(x_t,y_t) \to \kappa_\delta(x,y) := \frac{2\lambda}{1-\delta}\bigl(e^{(1-\delta)(x\wedge y)} - 1\bigr) \tag{20.3}$$

For $\delta = 1$, corresponding to $\delta = \gamma$ in the non-rescaled model, let

$$\kappa_1(x,y) := 2\lambda(x \wedge y). \tag{20.3'}$$

Then $\kappa_t^*(x_t,y_t) \to \kappa_\delta(x,y)$ in this case also.

Theorem 9.1 thus applies to $G(t)$ conditioned on the process $(N(t))_{t\ge0}$, and we find (conditioned on $(N(t))_{t\ge0}$, and thus also unconditionally) that

$$\frac{C_1\bigl(G(t)\bigr)}{N(t)} \xrightarrow{\mathrm{p}} \rho(\kappa_\delta),$$

with $\kappa_\delta$ given by (20.3) and (20.3').

To study $\rho(\kappa_\delta)$ further, and in particular to investigate the threshold as we vary $\lambda$ keeping $\mu \ge 0$ fixed, we thus have to investigate the integral operator $T_{\kappa_\delta}$ with kernel $\kappa_\delta$ given by (20.3). The change of variables $x \to e^{-x}$ transforms $\mathcal{S}$ and $\mu$ to the standard setting $(0,1]$ with Lebesgue measure, and the kernel (20.3) becomes

$$\tilde\kappa_\delta(x,y) := \frac{2\lambda}{1-\delta}\Bigl((x\vee y)^{\delta-1} - 1\Bigr), \tag{20.4}$$

with $\tilde\kappa_1 := 2\lambda \ln(1/(x\vee y))$.

In the case $\delta = 0$, this is the same as (18.1); hence we have the same critical value $1/8$ (for $\lambda$) as for the CHKNS model and the same $\rho(\kappa)$ giving the size of the giant component; in particular, the phase transition has infinite exponent. (Indeed, with $\delta = 0$ the model is very similar to (Durrett's form of) the CHKNS model discussed in Section 18; now a geometric number of edges between random vertices is added at each step, rather than a Poisson number.) For $\delta > 0$, the kernel $\tilde\kappa_\delta$ is in $L^2((0,1]^2)$, so $T_{\tilde\kappa_\delta}$ is compact and its norm can be found by finding its eigenvalues. These eigenvalues can be found by solving the Sturm–Liouville differential equation $\alpha G''(x) = -2\lambda x^{\delta-2}G(x)$ with boundary values $G(0) = G'(1) = 0$, where $\alpha$ denotes the eigenvalue.

The general solution satisfying $G(0)=0$ is, up to a constant factor, given by

$$G'(x)=\Bigl(\frac{2\lambda}{\alpha\delta^2}x^\delta\Bigr)^{-(1/\delta-1)/2}J_{1/\delta-1}\Bigl(2\Bigl(\frac{2\lambda}{\alpha\delta^2}x^\delta\Bigr)^{1/2}\Bigr),$$

where $J_\nu$ is a Bessel function. The remaining boundary condition $G'(1)=0$ (which gives the formula in Turova [45, Corollary 4.1] and [46]) thus leads to $J_{1/\delta-1}\Bigl(\Bigl(\frac{8\lambda}{\alpha\delta^2}\Bigr)^{1/2}\Bigr)=0$, so if $z_\nu$ is the first positive zero of $J_\nu$, then

$$\|T_{\kappa_\delta}\|=\|T_{\tilde\kappa_\delta}\|=\alpha_1=\frac{8\lambda}{\delta^2 z_{1/\delta-1}^2}.$$

In other words, the critical value of $\lambda$ is $\lambda_{\mathrm{cr}}(\delta)=\delta^2 z_{1/\delta-1}^2/8$, as given by a related argument by Söderberg [41].

Theorem 11.1 applies only when $\delta>1/2$, but the eigenfunctions are continuous and bounded for every $\delta>0$, and we believe that the phase transition has exponent 1, and that (11.1) holds, for every $\delta>0$.

We can easily find the asymptotics of $\lambda_{\mathrm{cr}}(\delta)$ as $\delta\to0$ or $\infty$; see Turova [45]. If $\lambda,\delta\to\infty$ with $\lambda/\delta\to c>0$, then $\tilde\kappa_\delta(x,y)\to 2c$, pointwise and in $L^2((0,1]^2)$, and it follows that $\|T_{\tilde\kappa_\delta}-T_{2c}\|\to0$. Consequently, for large $\delta$, the graph is subcritical if $2c<1$ and supercritical if $2c>1$. In other words, $\lambda_{\mathrm{cr}}(\delta)/\delta\to1/2$ as $\delta\to\infty$. Similarly, if $\delta\searrow0$, then $\kappa_\delta\nearrow\kappa_0$ and it follows easily that $\|T_{\kappa_\delta}\|\to\|T_{\kappa_0}\|$, and thus $\lambda_{\mathrm{cr}}(\delta)\to\lambda_{\mathrm{cr}}(0)=1/8$.

## 21. A "quantum random graph"

Ioffe and Levit [28] have recently introduced and studied a new random graph model related to quantum theory. To construct the graph, start with $n$ cirles of length $\beta>0$, and break the circles into pieces using independent Poisson processes of intensity $\lambda$. Then connect every pair of circles using a Poisson process with intensity $1/n$ of links.

It is easily seen that this is an instance of our model, where $\mathcal{S}$ is the family of all intervals in a circle of length $\beta$, $\mu$ is a certain measure with total mass $e^{-\lambda\beta}+\lambda\beta$ , and

$$\kappa(I,J)=|I\cap J|,$$

see Janson [30]. More precisely, $\mu$ is invariant undewr rotations and can thus be described by the induced measure $\widehat\mu$ on $\widehat{\mathcal{S}}:=(0,\beta]$ for the lengths of the intervals; this measure has a point mass $(1+\lambda\beta)e^{-\lambda\beta}$ at $\beta$ (corresponding to the "interval" which is the entire circle), and is otherwise continuous with density $\beta\lambda^2e^{-\lambda x}$ for $x\in(0,\beta)$.

A calculation [30], using the symmetry and reducing to the integral operator defined by $xy/\beta$ on $(\widehat{\mathcal{S}},\widehat\mu)$, shows that

$$\|T_\kappa\|=\frac{2}{\lambda}\bigl(1-e^{-\lambda\beta}\bigr)-\beta e^{-\lambda\beta}.$$

Theorem 9.1 applies, and the condition $\|T_\kappa\| > 1$ yields the threshold found by Ioffe and Levit [28] for the existence of a giant component.

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